

1) la ecuación diferencial es $\dot{T}_1(t) = \dot{T}(t) = 4(0 - 2T(t) + T_2(t))$

o, lo que es lo mismo: $\dot{T}(t) + 8T(t) = 4T_2(t)$

a) Para $t \leq 0$ es $T_2(t) = 0 \Rightarrow \dot{T}(t) + 8T(t) = 0 \Rightarrow T(t) = C e^{-8t}$
 que, con la condición inicial $T(0) = 0 \Rightarrow C = 0 \Rightarrow \boxed{T(t) = 0 \forall t \leq 0}$

Para $t > 0$ es $T_2(t) = e^{-8t} \Rightarrow \dot{T}(t) + 8T(t) = 4e^{-8t}$

Solución general de la ec. homogénea $T_h(t) = C e^{-8t}$

Existe resonancia \rightarrow Ensayamos $T_p(t) = A t e^{-8t}$

$\dot{T}_p + 8T_p = 4e^{-8t} \Rightarrow A = 4 \Rightarrow T(t) = C e^{-8t} + 4t e^{-8t}$

de la condición de empalme $T(0^-) = T(0^+) \Rightarrow 0 = C + 0 \Rightarrow C = 0$

Así $\boxed{T(t) = 4t e^{-8t} \forall t > 0}$

En definitiva, $\boxed{T(t) = 4t e^{-8t} u_s(t)}$ \rightarrow función salto.

b) $\dot{T}(t) + 8T(t) = 4T_2(t) \xrightarrow{\mathcal{F}} (i\omega + 8)\hat{T}(\omega) = 4\hat{T}_2(\omega)$

$\hat{T}_2(\omega) = \int_{-\infty}^{\infty} T_2(t) e^{-i\omega t} dt = \int_0^{\infty} e^{-(8+i\omega)t} dt = \frac{1}{8+i\omega} \Rightarrow \hat{T}(\omega) = \frac{4}{(8+i\omega)^2}$

Sabemos que $\mathcal{F}^{-1}\left(\frac{1}{8+i\omega}\right) = e^{-8t} u_s(t)$ y utilizando $\mathcal{F}(t^n f(t)) = i^n \frac{d^n \hat{f}(\omega)}{d\omega^n}$

y que $\frac{-i}{(8+i\omega)^2} = \frac{d}{d\omega} \left(\frac{1}{8+i\omega}\right) \Rightarrow \boxed{T(t) = \mathcal{F}^{-1}(\hat{T}(\omega)) = 4t e^{-8t} u_s(t)}$

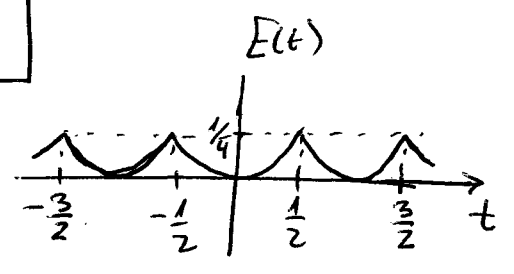
$\boxed{T(\infty) = \lim_{t \rightarrow \infty} T(t) = \lim_{t \rightarrow \infty} 4t e^{-8t} \stackrel{\text{L'Hopital}}{=} \lim_{t \rightarrow \infty} \frac{4}{8e^{8t}} = 0}$ Temperatura máxima

$T'(t) = 0 \Rightarrow (1-8t)e^{-8t} = 0 \Rightarrow \boxed{t_{\max} = \frac{1}{8}} \quad \boxed{T\left(\frac{1}{8}\right) = \frac{1}{2} e^{-1} \approx 0.184}$

2) $E(t) = t^2, t \in [-\frac{1}{2}, \frac{1}{2}], T=1, \omega_n = 2\pi n$

a) $C_n = \frac{1}{T} \int_{-\frac{1}{2}}^{\frac{1}{2}} t^2 e^{-i\omega_n t} dt = \left[\frac{u=t^2}{dv=e^{-i\omega_n t} dt} = \frac{t^2 e^{-i\omega_n t}}{-i\omega_n} \right]_{-\frac{1}{2}}^{\frac{1}{2}} + \int_{-\frac{1}{2}}^{\frac{1}{2}} 2t \frac{e^{-i\omega_n t}}{-i\omega_n} dt$
 $= \dots = 2 e^{-i\pi n} / (2\pi n)^2 = \frac{(-1)^n}{2\pi^2 n^2}$

$C_0 = \frac{1}{T} \int_{-\frac{1}{2}}^{\frac{1}{2}} t^2 dt = \left[\frac{t^3}{3} \right]_{-\frac{1}{2}}^{\frac{1}{2}} = \frac{1}{12}$



Al tratarse de una función par sólo tiene desarrollo en serie de cosenos:

$E(t) = \frac{A_0}{2} + \sum_{n=1}^{\infty} A_n \cos(\omega_n t) + \cancel{B_n \sin(\omega_n t)}$

$B_n = i(C_n - \bar{C}_n) = 0, A_n = C_n + \bar{C}_n = \frac{(-1)^n}{\pi^2 n^2}$

Con lo cual $E(t) = \frac{1}{12} + \sum_{n=1}^{\infty} \frac{(-1)^n}{\pi^2 n^2} \cos(2\pi n t) = \frac{1}{12} - \frac{\cos(2\pi t)}{\pi^2} + \dots$

b) $\ddot{I} + \dot{I} + I = \dot{E} \Rightarrow \hat{I}(\omega_n) (-\omega_n^2 + i\omega_n + 1) = i\omega_n \hat{E}(\omega_n)$

$V_L = L \dot{I}$
 Caída de tensión en la bobina $\Rightarrow H(\omega_n) = \frac{\hat{I}(\omega_n)}{\hat{E}(\omega_n)} = \frac{L i \omega_n \hat{I}(\omega_n)}{\hat{E}(\omega_n)} = \frac{L i \omega_n \hat{I}(\omega_n)}{L(-\omega_n^2 + i\omega_n + 1)}$
 (Note: The denominator in the original image is $1 - \omega_n^2 + i\omega_n$)

$|H(\omega_n)| = \left(\frac{\omega_n^4}{(1 - \omega_n^2)^2 + \omega_n^2} \right)^{1/2}$
 $\omega_n \rightarrow 0 \rightarrow 0$
 $\omega_n \rightarrow \infty \rightarrow 1$
 \Rightarrow filtro paso alta

Espectro de amplitud \uparrow

$$c) \hat{E}(t) = \frac{1}{L} - \frac{\cos(2\pi t)}{\pi^2}$$

$$\ddot{I} + R\dot{I} + \frac{1}{C}I = \hat{E}(t) = \frac{2}{\pi} \text{Sen}(2\pi t)$$

Solucion de la homogenea $I_h(t) = Ce^{\lambda t} \Rightarrow \lambda = \frac{-R \pm \sqrt{R^2 - \frac{4}{C}}}{2}$

Solo existe resonancia si $i\omega = \lambda \Rightarrow R=0, \frac{1}{C} = 2\pi \Rightarrow C = \frac{1}{4\pi^2}$

Asi, la ecuacion diferencial a resolver es:

$$\ddot{I} + 4\pi^2 I = \frac{2}{\pi} \text{Sen}(2\pi t) \text{ con cond. iniciales } I(0)=0 = \dot{I}(0)$$

$$I_h(t) = C_1 \cos(2\pi t) + C_2 \text{Sen}(2\pi t) \text{ (Solucion general de la ec. homog.)}$$

$$I_p(t) = At \cos(2\pi t) + Bt \text{Sen}(2\pi t) \text{ (Solucion particular)}$$

$$\ddot{I}_p + 4\pi^2 I_p = \frac{2}{\pi} \text{Sen}(2\pi t) \Rightarrow B=0, A = \frac{-1}{2\pi^2}$$

$$I(t) = I_h(t) + I_p(t) = C_1 \cos(2\pi t) + C_2 \text{Sen}(2\pi t) - \frac{t}{2\pi^2} \cos(2\pi t)$$

$$I(0) = \dot{I}(0) = 0 \Rightarrow C_1 = 0, C_2 = \frac{1}{4\pi^3} \Rightarrow$$

$$I(t) = \frac{1}{4\pi^3} \text{Sen}(2\pi t) - \frac{t}{2\pi^2} \cos(2\pi t)$$

3) a) $X_k - 2X_{k-1} + X_{k-2} = \delta_k$, $\lambda^2 - 2\lambda + 1 = 0 \Rightarrow \lambda = 1$ doble (4)

$$X_k = C_1 (1)^k + C_2 k (1)^k \quad \text{válida para } k \geq 1$$

$$X_0 = 2X_{-1} - X_{-2} + 1 = 1 = C_1 \left\{ \begin{array}{l} C_1 = 1 \\ C_2 = 1 \end{array} \right. \Rightarrow X_k = (1+k)u_k$$

$$X_{-1} = 0 = C_1 - C_2$$

b) $X_k - 2X_{k-1} + X_{k-2} = u_k$

$$X_p = A k^2 u_k \quad \text{con } A = \frac{1}{2}$$

$$X_k = \left(1 + \frac{3}{2}k + \frac{1}{2}k^2\right) u_k \quad (\text{Junio 2006})$$